

**A LINEARLY IMPLICIT FINITE ELEMENT METHOD
FOR A KLEIN–GORDON–SCHRÖDINGER-TYPE SYSTEM**

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ABSTRACT

Let $T > 0$, $\Omega := (x_A, x_B) \subset \mathbb{R}$ be a bounded interval and $D := [0, T] \times \bar{\Omega}$. Then, we consider the following initial- and boundary- value problem for a Klein–Gordon–Schrödinger (DKGS) system of partial differential equations: find functions $\psi : D \rightarrow \mathbb{C}$ and $\phi : D \rightarrow \mathbb{R}$ solving the following nonlinear system of equations

$$(0.1) \quad \psi_t = i \mu \psi_{xx} - \alpha \psi - i \phi \psi \quad \text{on } (0, T] \times \Omega,$$

$$(0.2) \quad \phi_{tt} = \phi_{xx} - \phi - \lambda \phi_t - \text{Re}(\psi_x) \quad \text{on } (0, T] \times \Omega,$$

and satisfying the conditions

$$(0.3) \quad \psi(t, x) = 0 \quad \text{and} \quad \phi(t, x) = 0 \quad \forall (t, x) \in [0, T] \times \partial\Omega,$$

$$(0.4) \quad \psi(0, x) = \psi_0(x) \quad \forall x \in \bar{\Omega},$$

$$(0.5) \quad \phi(0, x) = \phi_0(x) \quad \text{and} \quad \phi_t(0, x) = \phi_1(x) \quad \forall x \in \bar{\Omega},$$

where: $\mu > 0$, $\alpha > 0$, $\lambda > 0$ are known real numbers, and $\phi_0, \phi_1 : \bar{\Omega} \rightarrow \mathbb{R}$ and $\psi_0 : \bar{\Omega} \rightarrow \mathbb{C}$ are given functions. The DKGS system models the nonlinear interaction between high frequency electron waves and low frequency ion plasma waves in a homogeneous magnetic field. Such systems arise in the UHH plasma heating technique. The coupling of the two fluids is achieved through the (non-homogeneous) polarization drift and the induced current takes part in the collisional process of energy exchange. As a consequence, the nonlinearity differs from the one encountered in the Zakharov system [5]. For a detailed derivation of the model, along with the underlying assumptions, the reader is referred to [1] and [4]. Also, we refer to [1] for results on the existence and uniqueness of a solution of the problem. In the present work, we focus on the derivation of an efficient numerical method to approximate the solution of the initial and boundary value problem (0.1)-(0.5).

Let $r \in \mathbb{N}$ and $S_h^r \subset H_0^1(\Omega)$ be a finite element space consisting of functions which are piecewise polynomials of degree at most r over a partition of Ω in intervals with maximum mesh-length h . It is well-known that the following approximation property holds: there exists a constant $C_{A,r} > 0$ such that

$$(0.6) \quad \inf_{\chi \in S_h^r} \{ \|v - \chi\|_{L^2(\Omega)} + h \|v - \chi\|_{H^1(\Omega)} \} \leq C_{A,r} h^\ell \|v\|_{H^\ell(\Omega)} \quad \forall v \in H^\ell(\Omega) \cap H_0^1(\Omega), \quad \ell = 1, \dots, r+1.$$

Then, we define the discrete Laplacian operator $\Delta_h : H^1(\Omega) \rightarrow S_h^r$ by $(\Delta_h \varphi, \chi)_{L^2(\Omega)} = (\varphi', \chi')_{L^2(\Omega)}$ $\forall \chi \in S_h^r$, $\forall \varphi \in H^1(\Omega)$, and the L^2 -projection operator $P_h : L^2(D) \rightarrow S_h^r$ by $(P_h f - f, \chi)_{L^2(\Omega)} = 0$ $\forall \chi \in S_h^r$, $\forall f \in L^2(D)$. Also, we introduce elliptic projection operators $R_h^\psi, R_h^\phi : H_0^1(\Omega) \rightarrow S_h^r$ by $((R_h^\psi v - v)', \chi')_{L^2(\Omega)} = 0$ $\forall \chi \in S_h^r$, $\forall v \in H_0^1(\Omega)$, and $((R_h^\phi v - v)', \chi')_{L^2(\Omega)} + (R_h^\phi v - v, \chi)_{L^2(\Omega)} = 0$ $\forall \chi \in S_h^r$, $\forall v \in H_0^1(\Omega)$. Now, let $R_h = R_h^\psi$ or R_h^ϕ . The approximation property (0.6) (see [2]), yields

$$(0.7) \quad \|R_h v - \chi\|_{L^2(\Omega)} + h \|R_h v - \chi\|_{H^1(\Omega)} \leq C_{B,r} h^\ell \|v\|_{H^\ell(\Omega)} \quad \forall v \in H^\ell(\Omega) \cap H_0^1(\Omega), \quad \ell = 1, \dots, r+1,$$

Also, from [3] it holds that

$$(0.8) \quad \|R_h v - v\|_{L^\infty(\Omega)} \leq C_{r,r} h^\ell \|v\|_{W^{\ell,\infty}(\Omega)} \quad \forall v \in W^{\ell,\infty}(\Omega) \cap H_0^1(\Omega), \quad \ell = 1, \dots, r+1.$$

Let $N \in \mathbb{N}$, $k := \frac{T}{N}$ and $t^m := m k$ for $m = 0, \dots, N$. The proposed method constructs, for $m = 0, \dots, N$, an approximation $(\Psi_h^m, \Phi_h^m) \in S_h^r \times S_h^r$ of $(\psi(t^m, \cdot), \phi(t^m, \cdot))$ following the steps below:

Step 1: Set

$$(0.9) \quad \Psi_h^0 := R_h^\psi \psi_0 \quad \text{and} \quad \Phi_h^0 := R_h^\phi \phi_0.$$

Step 2: Set

$$(0.10) \quad \Phi_h^1 := R_h^\phi \left\{ \phi_0 + k \phi_1 + \frac{k^2}{2} [\phi_0'' - \phi_0 - \lambda \phi_1 - \text{Re}(\psi_0')] \right\}.$$

Step 3: Find $\Psi_h^1 \in S_h$ such that

$$(0.11) \quad \frac{\Psi_h^1 - \Psi_h^0}{k} = i \mu \Delta_h \left(\frac{\Psi_h^1 + \Psi_h^0}{2} \right) - \alpha \frac{\Psi_h^1 + \Psi_h^0}{2} - i P_h \left[\frac{\Phi_h^1 + \Phi_h^0}{2} \frac{\Psi_h^1 + \Psi_h^0}{2} \right].$$

Step 4: For $n = 1, \dots, N-1$, specify $(\Psi_h^{n+1}, \Phi_h^{n+1}) \in S_h \times S_h$ via the requirements

$$(0.12) \quad \frac{\Psi_h^{n+1} - \Psi_h^{n-1}}{2k} = i \mu \Delta_h \left(\frac{\Psi_h^{n+1} + \Psi_h^{n-1}}{2} \right) - \alpha \frac{\Psi_h^{n+1} + \Psi_h^{n-1}}{2} - i P_h \left[\Phi_h^n \frac{\Psi_h^{n+1} + \Psi_h^{n-1}}{2} \right]$$

and

$$(0.13) \quad \frac{\Phi_h^{n+1} - 2\Phi_h^n + \Phi_h^{n-1}}{k^2} = \Delta_h \left(\frac{\Phi_h^{n+1} + \Phi_h^{n-1}}{2} \right) - \lambda \frac{\Phi_h^{n+1} - \Phi_h^{n-1}}{2k} - \frac{\Phi_h^{n+1} + \Phi_h^{n-1}}{2} - P_h \left[\text{Re}((\Psi_h^n)') \right].$$

Analyzing the numerical method above, first we ensure that it is well-defined without mesh conditions, and then, we derive an a priori bound for the numerical approximations, in a discrete energy norm, described below

$$(0.14) \quad \max_{0 \leq m \leq N} \{ \|\Phi_h^m\|_1 + \|\Psi_h^m\|_1 \} + \max_{0 \leq n \leq N-1} \|\partial_k \Phi_h^n\|_0 \leq C \left[\|\partial_k \Phi_h^0\|_0 + \|\Phi_h^1\|_1 + \|\Phi_h^0\|_1 + \|\Psi_h^0\|_1 + \|\Psi_h^0\|_0^3 \right].$$

Finally, we combine (0.14), (0.6), (0.7) and (0.8), with a proper stability argument, to arrive at an optimal order error estimate in the L^2 and H^1 norms, i.e.,

$$\max_{0 \leq m \leq N} \left\{ \|\Psi_h^m - \psi^m\|_{H^1(\Omega)} + \|\Phi_h^m - \phi^m\|_{H^1(\Omega)} \right\} \leq C (k^2 + h^r)$$

and

$$\max_{0 \leq m \leq N} \left\{ \|\Psi_h^m - \psi^m\|_{L^2(\Omega)} + \|\Phi_h^m - \phi^m\|_{L^2(\Omega)} \right\} \leq C (k^2 + h^{r+1}).$$

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